# Towards nonlinear cell population model structured by molecular content

# Romain Yvinec

Inra Tours - Match Heidelberg

Joint work with M. Mackey, M. Tyran-Kaminska, A. Marciniak-Czochra





#### Stochasticity in Molecular biology, links with cell fate

Single cell model : Bursting and Division as Jump Processes

Nonlinear (macroscopic) population model Theoretical results Numerical results

#### Stochasticity in Molecular biology, links with cell fate

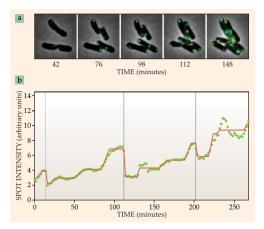
Single cell model : Bursting and Division as Jump Processes

Nonlinear (macroscopic) population model



# Stochasticity in molecular biology

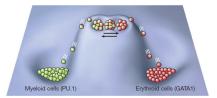
 Trajectories on single cells : bursting and repartition at division are major sources of randomness in gene expression.



[Golding et al. Cell 2005, Kondev Physics Today 2014] – action of the second s

# A typical example linking gene expression to cell fate

The antagonism between regulatory proteins (Transcription Factor) Gata-1/PU.1 in heamatopoietic progenitor



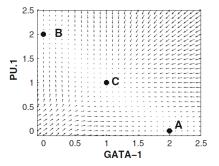
#### [Enver et al. Stem Cell 2009]

## Cell fate explained by a deterministic dynamical system

The antagonism Gata-1/PU1, modeled by ODE

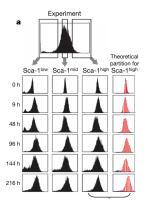
$$\frac{d[G]}{dt} = a_1 \frac{[G]^n}{\theta_{a1}^n + [G]^n} + b_1 \frac{\theta_{b1}^n}{\theta_{b1}^n + [P]^n} - k_1[G]$$
$$\frac{d[P]}{dt} = a_2 \frac{[P]^n}{\theta_{a2}^n + [P]^n} + b_2 \frac{\theta_{b2}^n}{\theta_{b2}^n + [G]^n} - k_2[P]$$

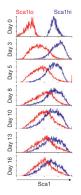
Cell fate "=" attractor of a dynamical system.



[Duff et al. JMB 2012]

# Reversibility of the gene expresion profile (and cell fate?) in *in vitro* cell culture experiment





[Chang et al. Nature Letters 08]

[Pina et al. Nature cell bio. 2012]

#### Stochasticity in Molecular biology, links with cell fate

#### Single cell model : Bursting and Division as Jump Processes

#### Nonlinear (macroscopic) population model



We define a pure-jump process  $(X(t))_{t\geq 0}$  on  $\mathbb{R}^*_+$  with two different transitions :

- Bursting at rate  $\lambda_b(x)$  and jump distribution  $\kappa_b(y, x) \mathbf{1}_{\{y>x\}} dy$
- Division at rate  $\lambda_d(x)$  and jump distribution  $\kappa_d(y, x) \mathbf{1}_{\{y < x\}} dy$

**Pathwise construction :** Let  $(U_n, V_n)_{n \ge 1}$  be i.i.d  $\propto \mathbb{U}(0, 1)$ ,

• Time step :  $T_n = T_{n-1} + (1/\lambda(X_{n-1})) \ln(1/U_{n-1})$ , where

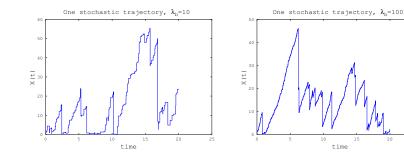
$$\lambda(x) = \lambda_b(x) + \lambda_d(x).$$

State step : X<sub>n</sub> = F<sup>-1</sup><sub>κ</sub>(V<sub>n</sub>, X<sub>n-1</sub>), where F<sub>κ</sub>(y, x) is the cum. dist. fonct. associated to

$$\kappa(y,x) = \underbrace{\frac{\lambda_b(x)}{\lambda_b(x) + \lambda_d(x)} \kappa_b(y,x) \mathbf{1}_{\{y > x\}}}_{\text{Bursting (gain)}} + \underbrace{\frac{\lambda_d(x)}{\lambda_b(x) + \lambda_d(x)} \kappa_d(y,x) \mathbf{1}_{\{y < x\}}}_{\text{Division (loss)}}.$$

$$\blacktriangleright X(t) = X_{n-1} \text{ for all } T_{n-1} \leq t < T_n.$$

# Example of sample paths



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This model is well-defined up to the explosion time,

$$T_{\infty} = \lim_{n \to \infty} T_n$$

#### Remark

Non-explosion + irreducibility + Existence of a unique invariant measure  $\Rightarrow$  ergodicity.

Lyapounov-fonction strategy (see [Meyn and Tweedie 93]) can provide sufficient criteria.

• An analogous study on the set of probability density  $(\int u = 1)$ .

$$\frac{\partial u(t,x)}{\partial t} = \underbrace{-\lambda_b(x)u(t,x) + \int_0^x \lambda_b(y)u(t,y)\kappa_b(x,y)dy}_{\text{Bursting (gain)}} \underbrace{-\lambda_d(x)u(t,x) + \int_x^\infty \lambda_d(y)u(t,y)\kappa_d(x,y)dy}_{\text{Division (loss)}}$$

This defines a semi-group P(t) on  $L^1$ . We will use

Theorem (Pichor and Rudnicki JM2A 2000) If P(t) is a stochastic semigroup :  $||P(t)u||_1 = ||u||_1$ , is partially integral, e.g. there exists  $t_0 > 0$  and p s.t.

$$\int_0^\infty \int_0^\infty p(x,y)\,dy\,dx>0 \quad \text{and} \quad P(t_0)u(x)\geq \int_0^\infty p(x,y)u(y)\,dy$$

and if P(t) possess a unique invariant density, then P(t) is asymptotically stable.

The Master equation may be rewritten as

$$\frac{du}{dt} = -\lambda u + K(\lambda u), \tag{1}$$

where

$$Kv(x) = \underbrace{\int_{0}^{x} \frac{\lambda_{b}(y)}{\lambda_{b}(y) + \lambda_{d}(y)} u(t, y)\kappa_{b}(x, y)dy}_{\text{Bursting (gain)}} + \underbrace{\int_{x}^{\infty} \frac{\lambda_{d}(y)}{\lambda_{b}(y) + \lambda_{d}(y)} u(t, y)\kappa_{d}(x, y)dy}_{\text{Division (loss)}}$$

If K has a strictly positive fixed point in  $L^1$ , then P(t) is stochastic ([Mackey et al. SIAM 13]). Note also that any stationary solution  $u^*$  of (1) must satisfy the flux condition

$$\underbrace{\int_{0}^{x} \left(\int_{x}^{\infty} \kappa_{b}(z,y) dz\right) \lambda_{b}(y) u^{*}(y) dy}_{\text{"from } x^{-} \text{ to } x^{+} \text{"}} = \underbrace{\int_{x}^{\infty} \left(\int_{0}^{x} \kappa_{d}(z,y) dz\right) \lambda_{d}(y) u^{*}(y) dy}_{\text{"from } x^{+} \text{ to } x^{-} \text{"}}$$

We consider the separable case

$$\kappa_b(x,y) = -rac{K_b'(x)}{K_b(y)}, \quad x > y, \quad \kappa_d(x,y) = rac{K_d'(x)}{K_d(y)}, \quad x < y.$$

where  $K_b(y) \to 0$  as  $y \to \infty$  and  $K(y) \to 0$  as  $y \to 0$ . We define

$$G(x) = \frac{K'_d(x)}{K_d(x)} - \frac{K'_b(x)}{K_b(x)}, \quad Q_b(x) = \int_x^{\overline{x}} \frac{\lambda_b(y)}{\lambda(y)} G(y) dy.$$

#### Theorem

Suppose that

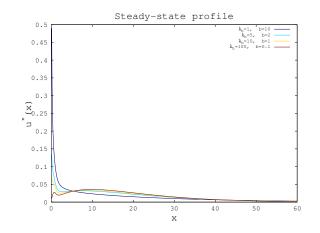
$$c_b := \int_0^\infty \frac{K_b(x)}{\lambda(x)} G(x) e^{-Q_b(x)} dx < \infty, \quad \int_0^\infty K_b(x) G(x) e^{-Q_b(x)} dx < \infty$$

Then the semigroup  $\{P(t)\}_{t\geq 0}$  is stochastic and is asymptotically stable, with

$$u_*(x) = \frac{1}{c_b} \frac{K_b(x)}{\lambda(x)} G(x) e^{-Q_b(x)}$$

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$$\frac{du^*}{dx} = \Big[ -\frac{\lambda'(x)}{\lambda(x)} + \frac{K_b'(x)}{K_b(x)} + \frac{G'(x)}{G(x)} + \frac{\lambda_b(x)}{\lambda(x)}G(x) \Big] u^*(x)$$



 $\mathcal{K}_b(x) = e^{-x/b}$ ,  $\lambda_b(x) = \lambda_b \frac{1+x^n}{\Lambda+x^n}$ ,  $\mathcal{K}_d(x) = x$ ,  $\lambda_d(x) = 1$ .

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This theorem can be used to show asymptotic convergence for "non-trivial" parameters function.

In particular, the growth-division model

$$\frac{\partial u(t,x)}{\partial t} + \underbrace{\frac{\partial g(x)u(t,x)}{\partial x}}_{\text{Continuous production}} = \underbrace{-\lambda_d(x)u(t,x) + \int_x^\infty \lambda_d(y)u(t,y)\frac{K_d'(x)}{K_d(y)}dy}_{\text{Division (loss)}},$$

converges for

$$\lambda_d(x) = \alpha x^{\beta-1} + x^{\beta+1}$$
$$g(x) = x^{\beta}$$
$$K_d(x) = x,$$
for  $0 \le \beta \le 1, \ 0 < \alpha < 1$ , towards
$$u_*(x) = \frac{K_d(x)}{\pi(x)} e^{-\int_{\bar{x}}^x \frac{\lambda_d(y)}{g(y)} dy}$$

but

$$\frac{\lambda_d}{g} \notin L^1_0$$

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cg(x)

**Absorbing probabilities / Mean waiting time :** We can also solve (analytically) the backward equation, Af(x) = A(x),

$$\mathcal{A}f(x) = \underbrace{\lambda_b(x) \left( \int_x^{\infty} (f(y) - f(x)) \kappa_b(y, x) dy \right)}_{\text{Bursting (gain)}} + \underbrace{\lambda_d(x) \left( \int_0^x (f(y) - f(x)) \kappa_d(y, x) dy \right)}_{\text{Division (loss)}}.$$

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$$\tau_z^+ := \inf\{t \ge 0, X_t \ge z\},\$$

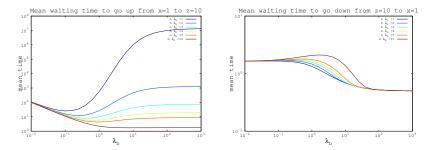
then

$$V_z^+(y) = \mathbb{E}_y[\tau_z^+]$$

is solution of

$$\begin{cases} \mathcal{A}V_z^+(y) = -1, \quad y < z, \\ V_z^+(y) = 0, \quad y \ge z. \end{cases}$$
(2)

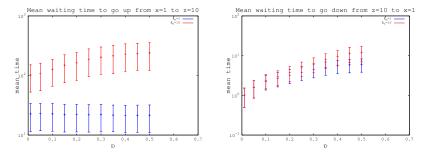
# The mean waiting time is non-monotonic with respect to the bursting property.



$$\lambda_d \equiv$$
 2,  $K_d(x) = x$ ,  $\lambda_b(x) \equiv \lambda_b$ ,  $K_b(x) = e^{-x/b}$ 

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# The mean waiting time is also affected by the asymmetry of the division.



$$\lambda_d(x) \equiv 2, \ \kappa_d(\cdot, x) = 0.5 \mathcal{N}(xp, xp(1-p)) + 0.5 \mathcal{N}(x(1-p), xp(1-p)), \ \kappa_b(x) = e^{-x/b}, \ b\lambda_b = 2$$

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#### Stochasticity in Molecular biology, links with cell fate

Single cell model : Bursting and Division as Jump Processes

Nonlinear (macroscopic) population model

Theoretical results Numerical results We wish to investigate (macroscopic) population models with nonlinear feedback on the division rate

$$\frac{\partial u(t,x)}{\partial t} = \underbrace{-\lambda_b(x)u(t,x) + \int_0^x \lambda_b(y)u(t,y)\kappa_b(x,y)dy}_{\text{Bursting}}$$

$$\underbrace{-\lambda_d(x, \mathbf{S})u(t,x) + 2\int_x^\infty \lambda_d(y, \mathbf{S})u(t,y)\kappa_d(x,y)dy}_{\text{Division}} - \underbrace{\mu(x)u(t,x)}_{\text{Cell death}}$$

with  $\kappa_d$  symmetric (total molecular content preserved at division) the feeback strength is given by

$$S(t) = \int_0^\infty \psi(x)u(t,x)dx, \quad \psi(x) = \mathbf{1}_{\{x \ge x_0\}}.$$

We will restrict to the case of *constant* division and death rates, so that

$$\frac{d}{dt}\left(\int_0^\infty u(t,x)dx\right) = (\lambda(S) - \mu)\int_0^\infty u(t,x)dx$$

If all cells participate to the regulation of the division rate  $(x_0 = 0)$ , we have immediately

#### Theorem

Let 
$$\kappa_b(x,y) = -\frac{\kappa_b'(x)}{\kappa_b(y)}$$
, and  $\kappa_d(x,y) = \frac{\kappa_d'(x)}{\kappa_d(y)}$ . We assume

$$c_b := \int_0^\infty \frac{K_b(x)}{\lambda(x)} G(x) e^{-Q_b(x)} dx < \infty, \quad \int_0^\infty K_b(x) G(x) e^{-Q_b(x)} dx < \infty$$

and that  $S \mapsto \lambda_d(S)$  is continuous monotonically decreasing, with  $\lambda_d(0) > \mu$  and  $\lim_{S \to \infty} \lambda_d(S) < \mu$ , then, for any initial density  $u_0$ , u(t, x) converges as  $t \to \infty$  in  $L^1$  towards

$$\lambda_d^{-1}(\mu)u^*.$$

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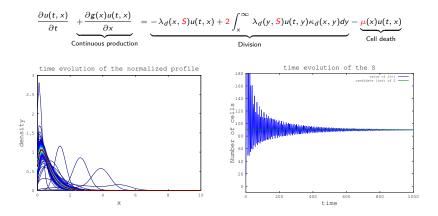
In the case  $x_0 > 0$ , we can not prove convergence towards a steady-state, and numerical results indicate the presence of oscillation through a Hopf-bifurcation.

#### Remark

We can however prove persistance results in certain cases

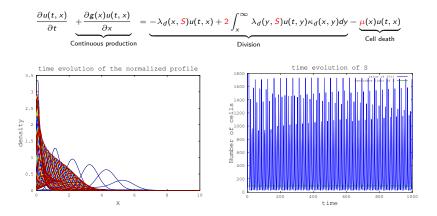
$$0 < \inf_{t \ge 0} \int_0^\infty u(t, x) dx \le \sup_{t \ge 0} \int_0^\infty u(t, x) dx < \infty$$

## Numerical results



$$\mu = 1$$
,  $\lambda_d(x, S) \equiv rac{10}{1+0.1*S}$ ,  $K_d(x) = x$ ,  $x_0 = 1$ ,  $g(x) \equiv 0.6$ 

## Numerical results indicate a Hopf bifurcation



$$\mu = 1$$
,  $\lambda_d(x, S) \equiv \frac{10}{1+0.1*S}$ ,  $K_d(x) = x$ ,  $x_0 = 1$ ,  $g(x) \equiv 0.5$ 

# The bursting property shifts the Hopf bifurcation

Table : +=Asymptotic convergence towards steady state - = oscillation

## The asymmetry at division also shifts the Hopf bifurcation

$$\frac{\frac{\partial u(t,x)}{\partial t}}{\frac{\partial t}{\partial t}} \underbrace{\underbrace{+\frac{\partial g(x)u(t,x)}{\partial x}}_{Continuous production}}_{Continuous production} = \underbrace{-\lambda_d(x,S)u(t,x) + 2\int_x^{\infty} \lambda_d(y,S)u(t,y)\kappa_d(x,y)dy}_{Division} - \underbrace{\frac{\mu(x)u(t,x)}{Cell \, death}}_{Cell \, death}$$
with  $\mu = 1$ ,  $\lambda_d(x,S) \equiv \frac{10}{1+0.1*S}$ ,  $\kappa_d(\cdot,x) = 0.5\mathcal{N}(xp,xp(1-p)) + 0.5\mathcal{N}(x(1-p),xp(1-p))$ ,  $x_0 = 1$ ,  $g(x) \equiv g$ 

$g \setminus p$	0.5	0.4	0.2	0.1	0.01
0.7	-	+	+	+	+
0.6	-	-	+	+	+
0.5	-	-	-	-	+

Table : +=Asymptotic convergence towards steady state - = oscillation

Upon an assumption of **separable bursting and division kernel**, we found a complete characterisation of the single cell model :

- Criteria for convergence towards steady-state, and analytical solution (and bifurcation)
- Mean waiting time to reach a given level

Such study can be used to infer the **burst rate** and/or **division rate** in a dividing cell population.

While looking at the nonlinear population model, the bursting properties and division mechanism are shown to have a profound impact on homeostasis that will be further investigated.

#### Thank you for your attention !