Towards nonlinear cell population model structured by molecular content

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Outline

Bursting and Division in gene expression models

Stochasticity in Molecular biology Bursting and Division as Jump Processes

Nonlinear population model

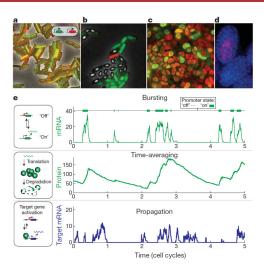
Theoretical results
Numerical results

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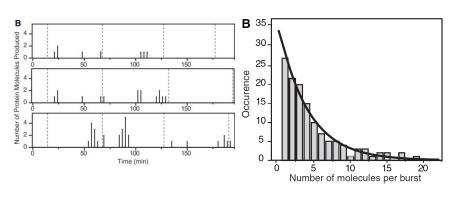


[Eldar and Elowitz Nature 2010]



Much more accurate measurements

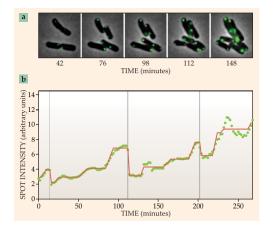
The bursting event are well characterized



[Yu et al. Science 06]

Much more accurate measurements

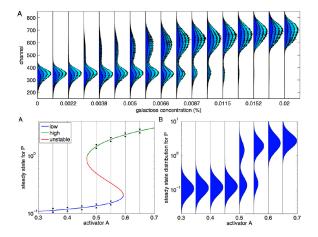
Trajectories can be analyzed on single cells.



[Golding et al. Cell 2005, Kondev Physics Today 2014]

Much more accurate measurements

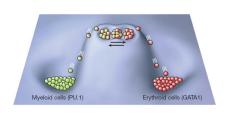
Bifurcation can be studied on probability distributions.



[Song et al. Plos CB 2010, Mackey et al. JTB 2011, SIAM 2013]

A typical example linking gene expression to cell fate

The antagonism Gata-1/PU.1 in heamatopoietic progenitor

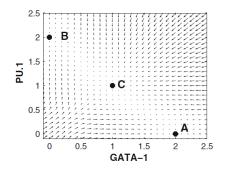


[Enver et al. Stem Cell 2009]

A typical example linking gene expression to cell fate

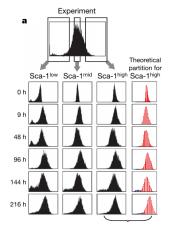
The antagonism Gata-1/PU1, modeled by ODE

$$\begin{split} \frac{d[G]}{dt} &= a_1 \frac{[G]^n}{\theta_{a_1}^n + [G]^n} + b_1 \frac{\theta_{b_1}^n}{\theta_{b_1}^n + [P]^n} - k_1[G] \\ \frac{d[P]}{dt} &= a_2 \frac{[P]^n}{\theta_{a_2}^n + [P]^n} + b_2 \frac{\theta_{b_2}^n}{\theta_{b_2}^n + [G]^n} - k_2[P] \end{split}$$



[Duff et al. JMB 2012]

A typical example linking gene expression to cell fate



Sca1hi Day 3 Day 5 Day 8 Day 13 Sca1

[Chang et al. Nature Letters 08]

[Pina et al. Nature cell bio. 2012]



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Numerical results

We define a pure-jump process $(X(t))_{t\geq 0}$ on \mathbb{R}_+^* with two different transitions :

- ▶ Bursting at rate $\lambda_b(x)$ and jump distribution $\kappa_b(y,x)\mathbf{1}_{\{y>x\}}dy$
- ▶ Division at rate $\lambda_d(x)$ and jump distribution $\kappa_d(y,x)\mathbf{1}_{\{y< x\}}dy$

Pathwise construction with the sequence $(U_n, V_n)_{n\geq 1}$, of i.i.d uniform random variable on (0,1)

 $ightharpoonup T_n = T_{n-1} - (1/\lambda(X_{n-1})) \ln(U_{n-1}), \text{ where}$

$$\lambda(x) = \lambda_b(x) + \lambda_d(x).$$

► $X_n = F_K^{-1}(V_n, X_{n-1})$, where $F_K(y, x)$ is the cum. dist. fonct. associated to

$$K(y,x) = \frac{\lambda_b(x)}{\lambda_b(x) + \lambda_d(x)} \kappa_b(y,x) \mathbf{1}_{\{y > x\}} + \frac{\lambda_d(x)}{\lambda_b(x) + \lambda_d(x)} \kappa_d(y,x) \mathbf{1}_{\{y < x\}}.$$

 $X(t) = X_{n-1}$ for all $T_{n-1} \le t < T_n$.



This model is well-defined up to the explosion time,

$$T_{\infty} = \lim_{n \to \infty} T_n$$

A well-known sufficient condition for non-explosion ($T_{\infty}=\infty$) is given by

$$\sum_{n>0} \frac{1}{\lambda_b(X_n) + \lambda_d(X_n)} = \infty.$$

In particular, this is the case for bounded jump rate.

Another criteria is provided by Lyapounov-fonction strategy (see [Mevn and Tweedie 93]). Let \mathcal{A} be the generator of $(X(t))_{t\geq 0}$,

$$\mathcal{A}f(x) = \lambda_b(x) \Big(\int_x^\infty (f(y) - f(x)) \kappa_b(y, x) dy \Big)$$
$$+ \lambda_d(x) \Big(\int_0^x (f(y) - f(x)) \kappa_d(y, x) dy \Big).$$

If there exists c > 0, V a positive measurable function s.t $V(x) \to \infty$ when $x \to 0$ and $x \to \infty$, $V \in \mathcal{D}(A)$ and

$$AV(x) \le cV(x), \quad x > 0,$$

then $(X(t))_{t>0}$ is non-explosif.

The fonction $V(x) = x^{-\gamma} \mathbf{1}_{\{x<1\}} + x^{\alpha} \mathbf{1}_{\{x>1\}}$ is suitable if there exists A, B, β, δ ,

$$ightharpoonup \overline{\kappa}_b(y,x) = \int_y^\infty \kappa_b(z,x) dz \le c(x/y)^\beta, \ \beta > \alpha$$

$$ightharpoonup \overline{\kappa}_d(y,x) = \int_0^y \kappa_d(z,x) dz \le c(y/x)^\delta, \ \delta > \gamma$$

 $\blacktriangleright \lambda_d(x) < A\lambda_b(x) + B \text{ as } x \to 0 \text{ and}$

$$\lim_{x\to 0}\lambda_b(x)x^{\delta}\int_x^1 y^{-\delta}\kappa_b(y,x)dy<\infty$$

 $\blacktriangleright \lambda_b(x) < A\lambda_d(x) + B \text{ as } x \to \infty \text{ and}$

$$\lim_{x\to\infty}\lambda_d(x)x^{-\alpha}\int_1^x y^\alpha\kappa_d(y,x)dy<\infty$$

Remark

"Similar" condition holds for ergodicity.

Remark

Non-explosion + irreductibility + Existence of a unique invariant $measure \Rightarrow ergodicity.$

▶ An analogous study on the set of probability density $(\int u = 1)$.

$$\begin{split} \frac{\partial u(t,x)}{\partial t} &= -\lambda_b(x)u(t,x) + \int_0^x \lambda_b(y)u(t,y)\kappa_b(x,y)dy \\ &- \lambda_d(x)u(t,x) + \int_x^\infty \lambda_d(y)u(t,y)\kappa_d(x,y)dy \end{split}$$

This defines a semi-group P(t) on L^1 . We will use

Theorem (Pichor and Rudnicki JM2A 2000)

If P(t)

- is a stochastic semigroup : $||P(t)u||_1 = ||u||_1$,
- ▶ is partially integral : there exists $t_0 > 0$ and p s.t.

$$\int_0^\infty \int_0^\infty p(x,y) \, dy \, dx > 0 \quad \text{and} \quad P(t_0) u(x) \ge \int_0^\infty p(x,y) u(y) \, dy$$

and possess a unique invariant density,

then P(t) is asymptotically stable.



The Master equation may be rewritten as

$$\frac{du}{dt} = -\lambda u + K(\lambda u),\tag{1}$$

where

$$Kv(x) = \int_0^x \frac{\lambda_b(y)}{\lambda_b(y) + \lambda_d(y)} u(t, y) \kappa_b(x, y) dy + \int_x^\infty \frac{\lambda_d(y)}{\lambda_b(y) + \lambda_d(y)} u(t, y) \kappa_d(x, y) dy$$

If K has a strictly positive fixed point in L^1 , then P(t) is stochastic ([Mackey et al. SIAM 13]). Note also that any stationary solution u^* of (1) must satisfy the flux condition

$$\int_0^x \overline{\kappa}_b(x,y) \lambda_b(y) u^*(y) dy = \int_x^\infty \overline{\kappa}_d(x,y) \lambda_d(y) u^*(y) dy$$

We consider the separable case

$$\kappa_b(x,y) = -\frac{K_b'(x)}{K_b(y)}, \quad x > y, \quad \kappa_d(x,y) = \frac{K_d'(x)}{K_d(y)}, \quad x < y.$$

where $K_b(y) \to 0$ as $y \to \infty$ and $K(y) \to 0$ as $y \to 0$. We define

$$G(x) = \frac{K_d'(x)}{K_d(x)} - \frac{K_b'(x)}{K_b(x)}, \quad Q_b(x) = \int_x^{\overline{x}} \frac{\lambda_b(y)}{\lambda(y)} G(y) dy.$$

Theorem

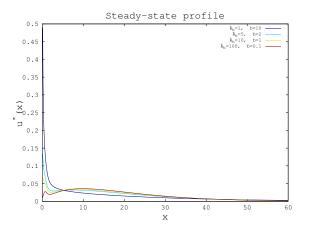
Suppose that

$$c_b := \int_0^\infty \frac{K_b(x)}{\lambda(x)} G(x) e^{-Q_b(x)} dx < \infty, \quad \int_0^\infty K_b(x) G(x) e^{-Q_b(x)} dx < \infty$$

Then the semigroup $\{P(t)\}_{t\geq 0}$ is stochastic and is asymptotically stable, with

$$u_*(x) = \frac{1}{c_b} \frac{K_b(x)}{\lambda(x)} G(x) e^{-Q_b(x)}$$

$$\frac{du^*}{dx} = \left[-\frac{\lambda'(x)}{\lambda(x)} + \frac{K_b'(x)}{K_b(x)} + \frac{G'(x)}{G(x)} + \frac{\lambda_b(x)}{\lambda(x)}G(x) \right] u^*(x)$$



$$K_b(x) = e^{-x/b}$$
, $\lambda_b(x) = \lambda_b \frac{1+x^n}{\Lambda+x^n}$, $K_d(x) = x$, $\lambda_d(x) = 1$.

▶ This theorem can be used to show asymptotic convergence for "non-trivial" parameters function.

In particular, the growth-division model

$$\frac{\partial u(t,x)}{\partial t} + \frac{\partial g(x)u(t,x)}{\partial x} = -\lambda_d(x)u(t,x) + \int_x^\infty \lambda_d(y)u(t,y)\frac{K'_d(x)}{K_d(y)}dy,$$

converges for

$$\lambda_d(x) = \alpha x^{\beta - 1} + x^{\beta + 1}$$
$$g(x) = x^{\beta}$$
$$K_d(x) = x,$$

for $0 < \beta < 1$. $0 < \alpha < 1$. towards

$$u_*(x) = \frac{K_d(x)}{cg(x)} e^{-\int_{\bar{x}}^x \frac{\lambda_d(y)}{g(y)} dy},$$

but

$$\frac{\lambda_d}{g} \notin L_0^1$$



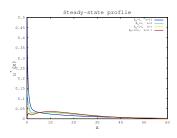
$$\tau_{u,z} := \inf\{t \geq 0, X_t \geq z\},\$$

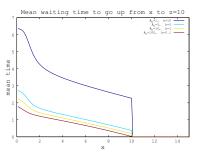
then

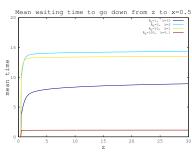
$$V_{u,z}(y) = \mathbb{E}_y[\tau_{u,z}]$$

is solution of

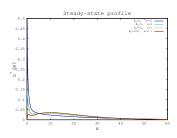
$$\begin{cases} A V_{u,z}(y) = -1, & y < z, \\ V_{u,z}(y) = 0, & y \ge z. \end{cases}$$
 (2)

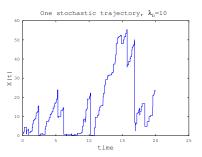


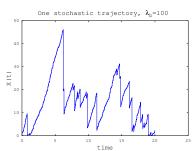




$$K_b(x) = e^{-x/b}$$
, $\lambda_b(x) = \lambda_b \frac{1+x^n}{\Lambda+x^n}$, $K_d(x) = x$, $\lambda_d(x) \equiv 1$.



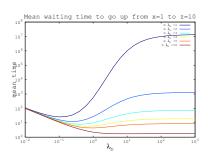


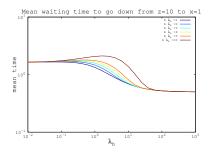


$$K_b(x) = e^{-x/b}$$
, $\lambda_b(x) = \lambda_b \frac{1+x^n}{\Lambda+x^n}$, $K_d(x) = x$, $\lambda_d(x) \equiv 1$.



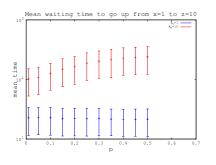
The mean waiting time is non-monotonic with respect to the bursting property.

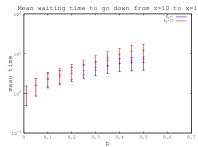




$$\lambda_d \equiv 2$$
, $K_d(x) = x$, $\lambda_b(x) \equiv \lambda_b$, $K_b(x) = e^{-x/b}$

The mean waiting time is also affected by the asymmetry of the division.





$$\lambda_d(x) \equiv 2,$$
 $K_d(x) = 0.5 \mathcal{N}(xp, xp(1-p)) + 0.5 \mathcal{N}(x(1-p), xp(1-p)),$
 $K_b(x) = e^{-x/b}, \ b\lambda_b = 2$

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Nonlinear population model

Theoretical results

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We wish to investigate (macroscopic) population models with nonlinear feedback on the division rate

$$\begin{split} \frac{\partial u(t,x)}{\partial t} &= -\lambda_b(x)u(t,x) + \int_0^x \lambda_b(y)u(t,y)\kappa_b(x,y)dy \\ &- \lambda_d(x,S)u(t,x) + 2\int_x^\infty \lambda_d(y,S)u(t,y)\kappa_d(x,y)dy - \mu(x)u(t,x) \end{split}$$

with κ_d symmetric (total molecular content preserved at division) the feeback strenght is given by

$$S(t) = \int_0^\infty \psi(x) u(t,x) dx, \quad \psi(x) = \mathbf{1}_{\{x \ge x_0\}}.$$

We will restrict to the case of *constant* division and death rates, so that

$$\frac{d}{dt}\Big(\int_0^\infty u(t,x)dx\Big) = (\lambda(S) - \mu)\int_0^\infty u(t,x)dx$$



If all cells participate to the regulation of the division rate $(x_0 = 0)$, we have immediately

Theorem

Let
$$\kappa_b(x,y) = -\frac{\kappa_b'(x)}{\kappa_b(y)}$$
, and $\kappa_d(x,y) = \frac{\kappa_d'(x)}{\kappa_d(y)}$. We assume

$$c_b := \int_0^\infty \frac{K_b(x)}{\lambda(x)} G(x) e^{-Q_b(x)} dx < \infty, \quad \int_0^\infty K_b(x) G(x) e^{-Q_b(x)} dx < \infty$$

and that $S \mapsto \lambda_d(S)$ is continuous monotonically decreasing, with $\lambda_d(0) > \mu$ and $\lim_{S \to \infty} \lambda_d(S) < \mu$, then, for any initial density u_0 , u(t,x) converges as $t \to \infty$ in L^1 towards

$$\lambda_d^{-1}(\mu)u^*$$
.



In the case $x_0 > 0$, we can only prove a persistance result for the equation

$$\frac{\partial u(t,x)}{\partial t} + \frac{\partial g(x)u(t,x)}{\partial x} = -\lambda_d(S)u(t,x) + 2\int_x^\infty \lambda_d(S)u(t,y)\kappa_d(x,y)dy - \mu u(t,x)$$

Theorem

With g smooth, bounded and bounded away from 0, starting with a positive $u_0 \in L^1$, we have

$$0 < \inf_{t \ge 0} \int_0^\infty u(t, x) dx \le \sup_{t \ge 0} \int_0^\infty u(t, x) dx < \infty$$
$$0 < \inf_{t \ge 0} S(t) \le \sup_{t \ge 0} S(t) < \infty$$

Démonstration.

We define $v(t,x):=e^{\int_0^t (\mu-\lambda_d(S(s)))ds}u(t,x)$, so that

$$\frac{\partial v(t,x)}{\partial t} + \frac{\partial g(x)v(t,x)}{\partial x} = -2\lambda_d(S)v(t,x) + 2\lambda_d(S)\int_x^\infty v(t,y)\kappa_d(x,y)dy$$

We use a coupling strategy to show that

$$\int_{x_0}^{\infty} v(t,x)dx \ge c(1+\varepsilon(t))$$

with $\varepsilon(t) \to 0$ (at exponential speed). For this, we use the coupling

$$Af(x,y) = g(x)f'(x) + g(y)f'(y)$$

$$+ 2\lambda_d(S(t)) \left(\int_0^1 (f(xz,yz) - f(x,y))dz \right)$$

$$+ 2(\|\lambda_d\|_{\infty} - \lambda_d(S(t))) \left(\int_0^1 (f(xz,y) - f(x,y))dz \right).$$

Then, $\int_{x_0}^{\infty} v(t,x) dx \ge \int_{x_0}^{\infty} w(t,x) dx$ where

$$\frac{\partial w(t,x)}{\partial t} + \frac{\partial g(x)w(t,x)}{\partial x} = -2\|\lambda_d\|_{\infty}w(t,x) + 2\|\lambda_d\|_{\infty}\int_{x}^{\infty}w(t,y)\kappa_d(x,y)dy$$

which converges as $t \to \infty$ due to hypotheses on g, κ_d .



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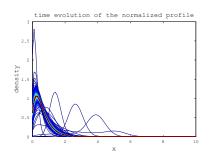
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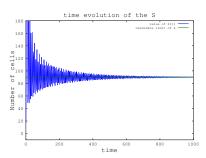
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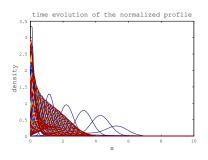


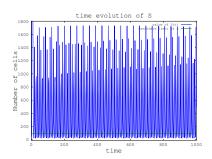


$$\mu = 1$$
, $\lambda_d(x, S) \equiv \frac{10}{1 + 0.1 * S}$, $K_d(x) = x$, $x_0 = 1$, $g(x) \equiv 0.6$



Numerical results indicate a Hopf bifurcation





$$\mu = 1$$
, $\lambda_d(x, S) \equiv \frac{10}{1 + 0.1 * S}$, $K_d(x) = x$, $x_0 = 1$, $g(x) \equiv 0.5$



The bursting property shifts the Hopf bifurcation : with $\mu=1$, $\lambda_d(x,S)\equiv \frac{10}{1+0.1*S}$, $K_d(x)=x$, $x_0=1$, $K_b(x)=e^{-x/b}$, $\lambda_b(x)\equiv \lambda_b$

$b\lambda_b \backslash \lambda_b$	100	10	1	0.1
0.6	+	+	+	+
0.5	-	+	+	+
0.4	-	-	+	+
0.1	-	-	-	+

 $\begin{tabular}{ll} \textbf{Table}: += & A symptotic convergence towards steady state -= oscillation \\ \end{tabular}$



The asymmetry at division also shifts the Hopf bifurcation: with $\mu = 1, \ \lambda_d(x, S) \equiv \frac{10}{1 + 0.1 * S}$ $\kappa_d(\cdot, x) = 0.5 \mathcal{N}(xp, xp(1-p)) + 0.5 \mathcal{N}(x(1-p), xp(1-p)),$ $x_0 = 1, g(x) \equiv g$

g\p	0.5	0.4	0.2	0.1	0.01
0.7	-	+	+	+	+
0.6	-	-	+	+	+
0.5	-	-	-	-	+

Table: +=Asymptotic convergence towards steady state -= oscillation

Vielen Dank!